Multibody Dynamics

Constraint Relaxation Method: Meaning of Lagrange Multipliers

O Previously, it was noted that if a dynamic system is described using "n" generalized coordinates q_k (k = 1,...,n), and if the system is subjected to "m" independent holonomic and/or nonholonomic constraint equations of the form

$$\left| \sum_{k=1}^{n} a_{jk} \dot{q}_{k} + a_{j0} \right| = 0 \qquad (j = 1, ..., m)$$
 (1)

tthe equations of motion of the system can be found by using one of the following two forms of Lagrange's equations with Lagrange multipliers.

$$\left| \frac{d}{dt} \left(\frac{\partial K}{\partial \dot{q}_k} \right) - \frac{\partial K}{\partial q_k} \right| = F_{q_k} + \sum_{j=1}^m \lambda_j a_{jk}$$
 (k = 1,...,n)

or

$$\left| \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_k} \right) - \frac{\partial L}{\partial q_k} \right| = \left(F_{q_k} \right)_{nc} + \sum_{j=1}^m \lambda_j a_{jk}$$
 (k = 1,...,n)

Here,

 $\succ K$: *kinetic energy* of the system

 $\succ F_{q_k}$: **generalized force** associated with the generalized coordinate q_k

> L : Lagrangian of the system

 $\triangleright V$: potential energy function for the conservative forces and torques

 $\triangleright (F_{q_k})_{nc}$: generalized force associated with q_k for only the nonconservative forces/torques

 $\triangleright \lambda_i$: Lagrange multiplier associated with the j^{th} constraint equation

 \triangleright a_{jk} : coefficients from the constraint equations (j=1,...,m; k=1,...,n)

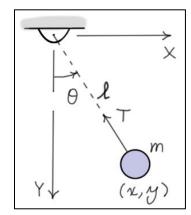
- Eqs. (1) and (2) or Eqs. (1) and (3) form a set of "n+m" differential/algebraic equations for the "n" generalized coordinates and the "m" Lagrange multipliers.
- Alternatively, some or all the constraints can be *relaxed* (or *removed*) and replaced with *force* and/or *torque* components that are required to *maintain* the *constraints*. Then, formulate the "n" Lagrange's equations in terms of the "n" generalized coordinates and the "m" constraint force (or torque) components.

Together with the constraint equations, this forms a set of "n+m" differential/ algebraic equations for the "n" generalized coordinates and the "m" constraint force and/or torque components. If all the constraints are relaxed, then Eqs. (3) can be written as

$$\left| \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_k} \right) - \frac{\partial L}{\partial q_k} \right| = \left(F_{q_k} \right)_{nc} + \left(F_{q_k} \right)_{\text{constraints}}$$
 (k = 1,...,n) (4)

Example: The Simple Pendulum

- \circ For the simple pendulum shown, $q_1 = x$ and $q_2 = y$ are used as the generalized coordinates, and the *length constraint* of the pendulum is *relaxed* in the formulation. Lagrange's equations can then be written in the form of Eq. (4).
- Here, $L = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) + mgy$, $(F_{q_k})_{nc} = 0$, and the contributions of the constraint force to the right sides of the equations are



$$\begin{aligned}
& \left[\left(F_{x} \right)_{\text{constraint}} = T \cdot \left(\partial y / \partial \dot{x} \right) = T \left(-\left(x / \ell \right) \dot{\underline{i}} - \left(y / \ell \right) \dot{\underline{j}} \right) \cdot \partial \left(\dot{x} \dot{\underline{i}} + \dot{y} \dot{\underline{j}} \right) / \partial \dot{x} = -T \left(x / \ell \right) \\
& \left[\left(F_{y} \right)_{\text{constraint}} = T \cdot \left(\partial y / \partial \dot{y} \right) = T \left(-\left(x / \ell \right) \dot{\underline{i}} - \left(y / \ell \right) \dot{\underline{j}} \right) \cdot \partial \left(\dot{x} \dot{\underline{i}} + \dot{y} \dot{\underline{j}} \right) / \partial \dot{y} = -T \left(y / \ell \right) \end{aligned} \tag{5}$$

$$(F_{y})_{\text{constraint}} = \underline{T} \cdot (\partial y / \partial \dot{y}) = T(-(x/\ell)\underline{i} - (y/\ell)\underline{j}) \cdot \partial (\dot{x}\underline{i} + \dot{y}\underline{j}) / \partial \dot{y} = -T(y/\ell)$$

$$(6)$$

Substituting into Lagrange's equations (4) and supplementing with the twice differentiated constraint equation gives the following equations of motion.

$$m\ddot{x} + \left(\frac{x}{\ell}\right)T = 0$$

$$m\ddot{y} - mg + \left(\frac{y}{\ell}\right)T = 0$$

$$x\ddot{x} + y\ddot{y} + \dot{x}^2 + \dot{y}^2 = 0$$
(7)

Using Lagrange multipliers, it was shown in previous notes that the equations for the pendulum could be written as

$$m\ddot{x} - \lambda x = 0$$

$$m\ddot{y} - mg - \lambda y = 0$$

$$x\ddot{x} + y\ddot{y} + \dot{x}^2 + \dot{y}^2 = 0$$
(8)

Comparing Eqs. (7) and (8), it is clear that the *Lagrange multiplier* $\lambda = -T/\ell$, the tension force per unit pendulum length.